



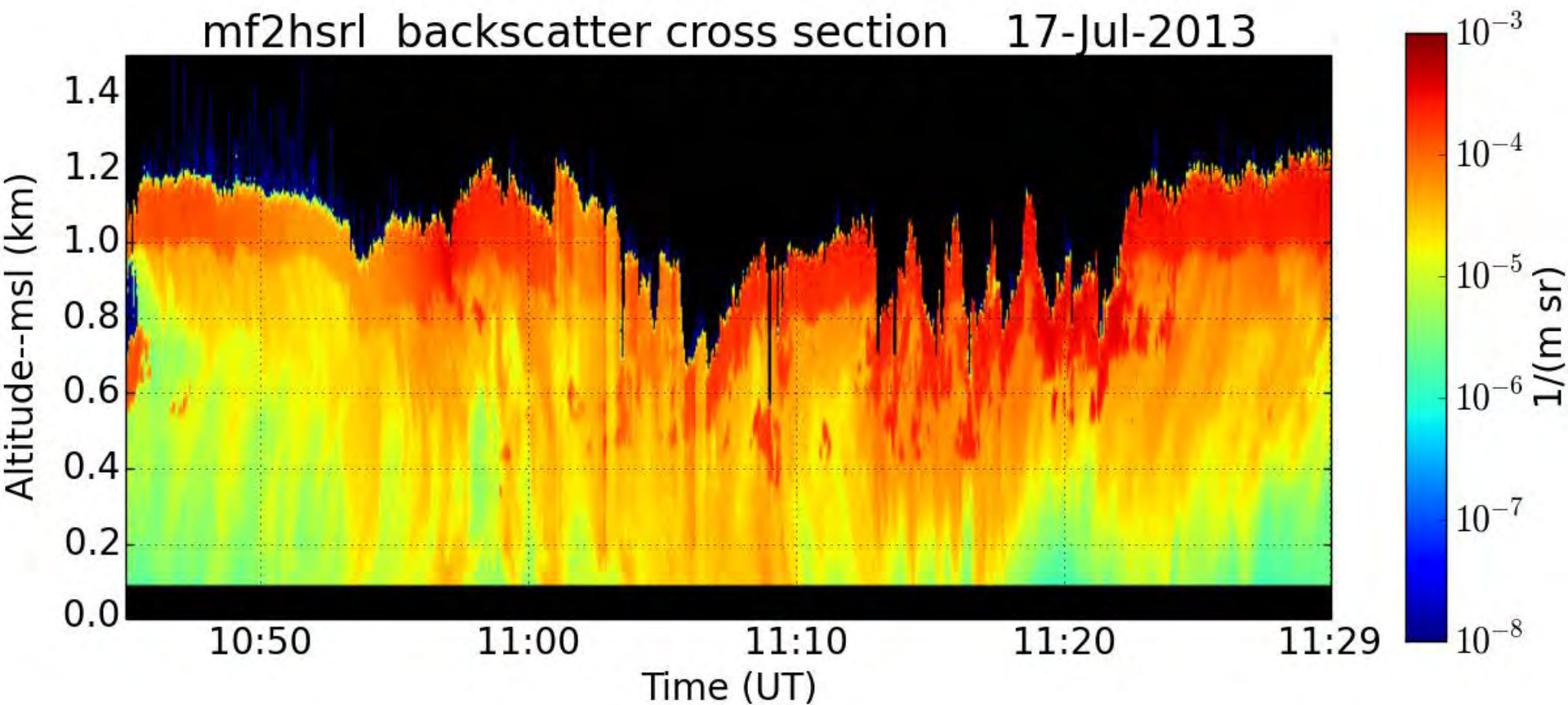
Drizzle measurements with the HSRL and the KAZR:
sensitivity to assumptions

Ed Eloranta

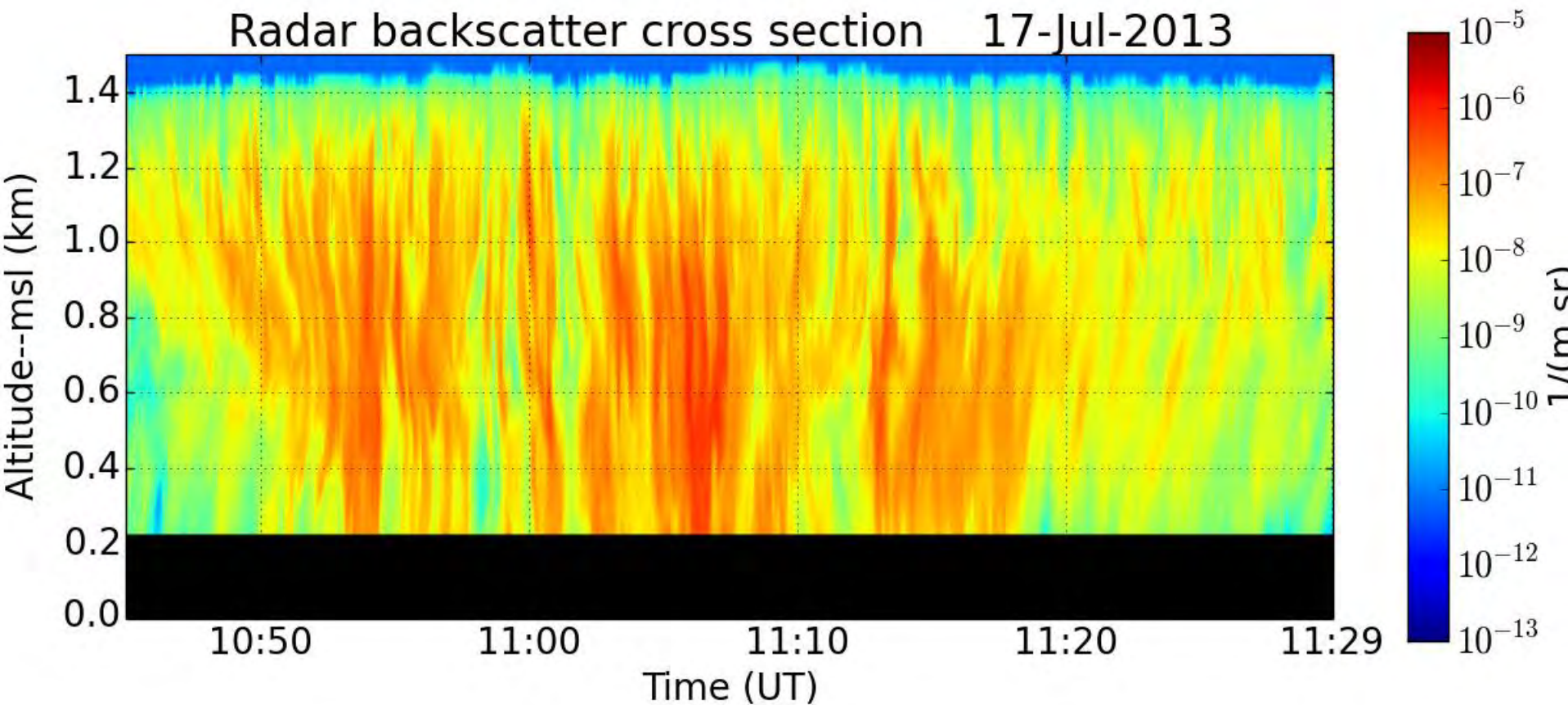
University of Wisconsin-Madison

<http://lidar.ssec.wisc.edu>

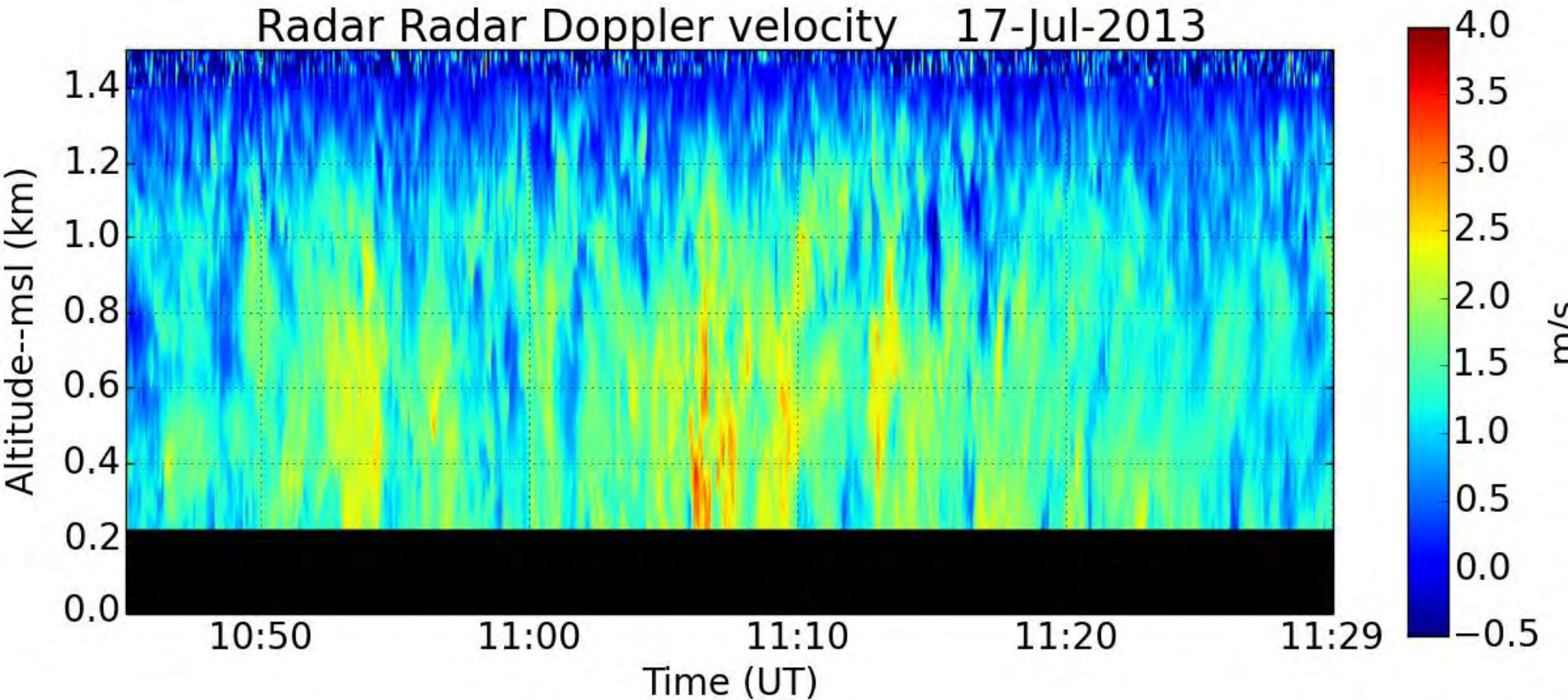
mf2hsrl backscatter cross section 17-Jul-2013



Radar backscatter cross section 17-Jul-2013



Radar Radar Doppler velocity 17-Jul-2013



HSRL-Radar particle size measurement

For droplets which are small compared to the radar wavelength but large compared to the lidar wavelength:

Radar scattering cross section = $\beta_{\text{radar}} \sim \langle V^2 \rangle \sim \langle D^6 \rangle$ ----- Rayleigh scattering
Lidar extinction cross section = $\beta_{\text{lidar}} \sim \langle A \rangle \sim \langle D^2 \rangle$ ----- Geometric optics

Where:

$\langle V^2 \rangle$ = average volume squared of the particles

$\langle A \rangle$ = average projected area of the particles

We can define:

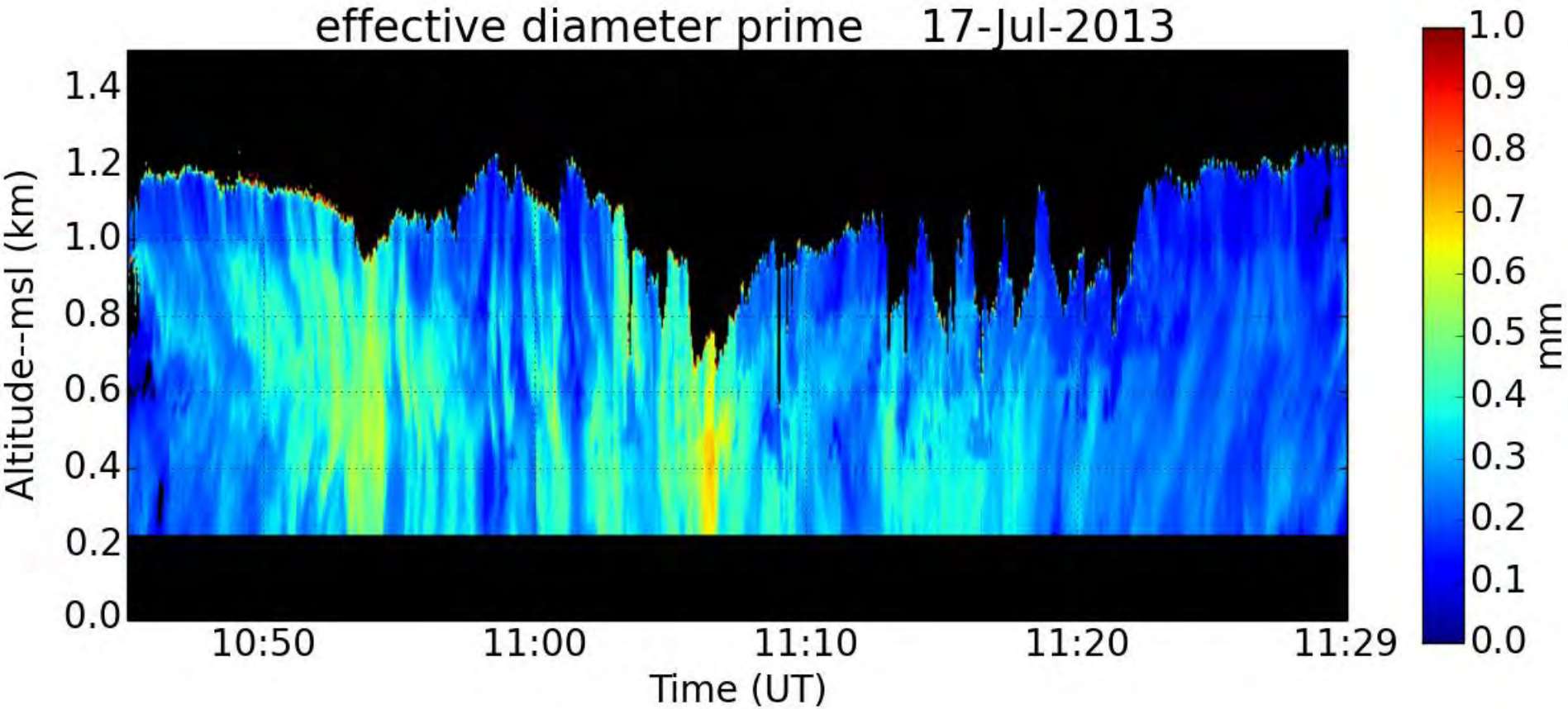
$$D'_{eff} = \left(\frac{\lambda_{\text{radar}}}{\pi} \right) \cdot \left(\frac{3 \cdot \beta_{\text{radar}}}{4 \cdot k_w^2 \cdot \beta_{\text{lidar}}} \right)^{\frac{1}{4}}$$

Where:

λ_{radar} = radar wavelength

k_w = dielectric constant of water

effective diameter prime 17-Jul-2013



Assuming a gamma distribution of particle diameters:

$$N = a \int_0^{\infty} D^{\alpha} \cdot \exp\left(-\frac{\alpha}{\gamma} \cdot \left(\frac{D}{D_m}\right)^{\gamma}\right) \cdot dD$$

Where:

N = number of particles

D = particle Diameter

D_m = mode diameter

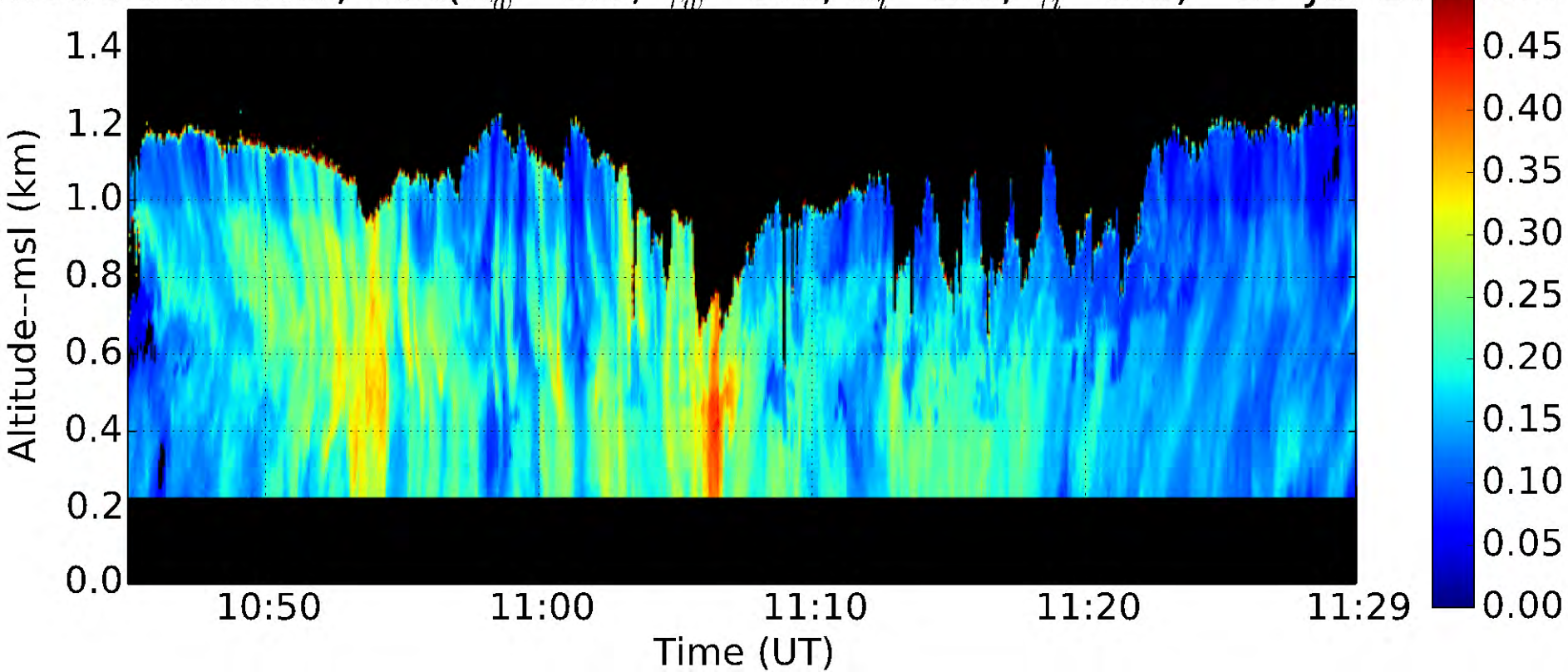
The effective diameter prime can be written as:

$$D'_{eff} = \frac{\int D^{\alpha+6} \cdot \exp\left(-\frac{\alpha}{\gamma} \cdot \frac{D}{D_m}^{\gamma}\right) \cdot dD}{\int D^{\alpha+2} \cdot \exp\left(-\frac{\alpha}{\gamma} \cdot \frac{D}{D_m}^{\gamma}\right) \cdot dD} = \frac{\frac{\alpha}{\gamma}^{\frac{1}{\gamma}} \cdot D_m \cdot \Gamma\left(\frac{\alpha+7}{\gamma}\right)}{\Gamma\left(\frac{\alpha+3}{\gamma}\right)}$$

Solving for the mode diameter, D_m :

$$D_m = D'_{eff} \cdot \left(\frac{\alpha}{\gamma}\right)^{\frac{1}{\gamma}} \cdot \left(\frac{\Gamma\left(\frac{\alpha+3}{\gamma}\right)}{\Gamma\left(\frac{\alpha+7}{\gamma}\right)}\right)^{\frac{1}{4}}$$

mode diameter, dist($\alpha_w = 1.0, \gamma_w = 3.0, \alpha_i = 1.0, \gamma_i = 1.0$) 17-Jul-2013 0.50



The radar weighted fall velocity, $\langle V_{rf} \rangle$:

$$\langle V_{rf} \rangle = \frac{\int V_f \cdot D^{\alpha+6} \cdot \exp\left(-\frac{\alpha}{\gamma} \cdot \frac{D}{D_m}^\gamma\right) \cdot dD}{\int D^{\alpha+6} \cdot \exp\left(-\frac{\alpha}{\gamma} \cdot \frac{D}{D_m}^\gamma\right) \cdot dD}$$

And the mass weighted fall velocity, $\langle V_{mf} \rangle$:

$$\langle V_{mf} \rangle = \frac{\int V_f \cdot D^{\alpha+3} \cdot \exp\left(-\frac{\alpha}{\gamma} \cdot \frac{D}{D_m}^\gamma\right) \cdot dD}{\int D^{\alpha+3} \cdot \exp\left(-\frac{\alpha}{\gamma} \cdot \frac{D}{D_m}^\gamma\right) \cdot dD}$$

Where the fall velocity, V_f , is computed from:

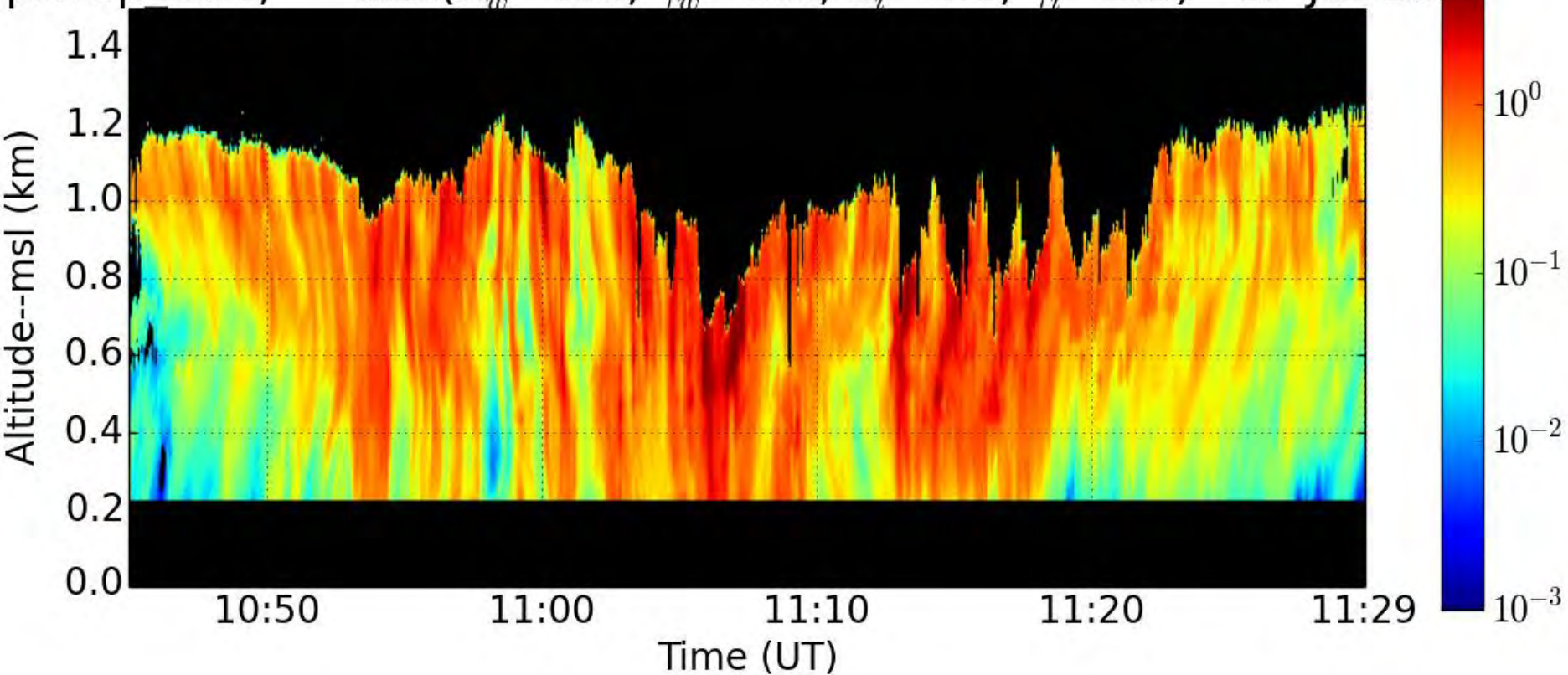
$$V_f = \frac{\eta}{\rho_{air} D} \left(\frac{\delta_0^2}{4} \left[\left(1 + C_1 X^{\frac{1}{2}}\right)^{\frac{1}{2}} - 1 \right]^2 - a_0 X^{b_0} \right)$$

Khovostyanov and Curry, JAS May 2005, Vol 62

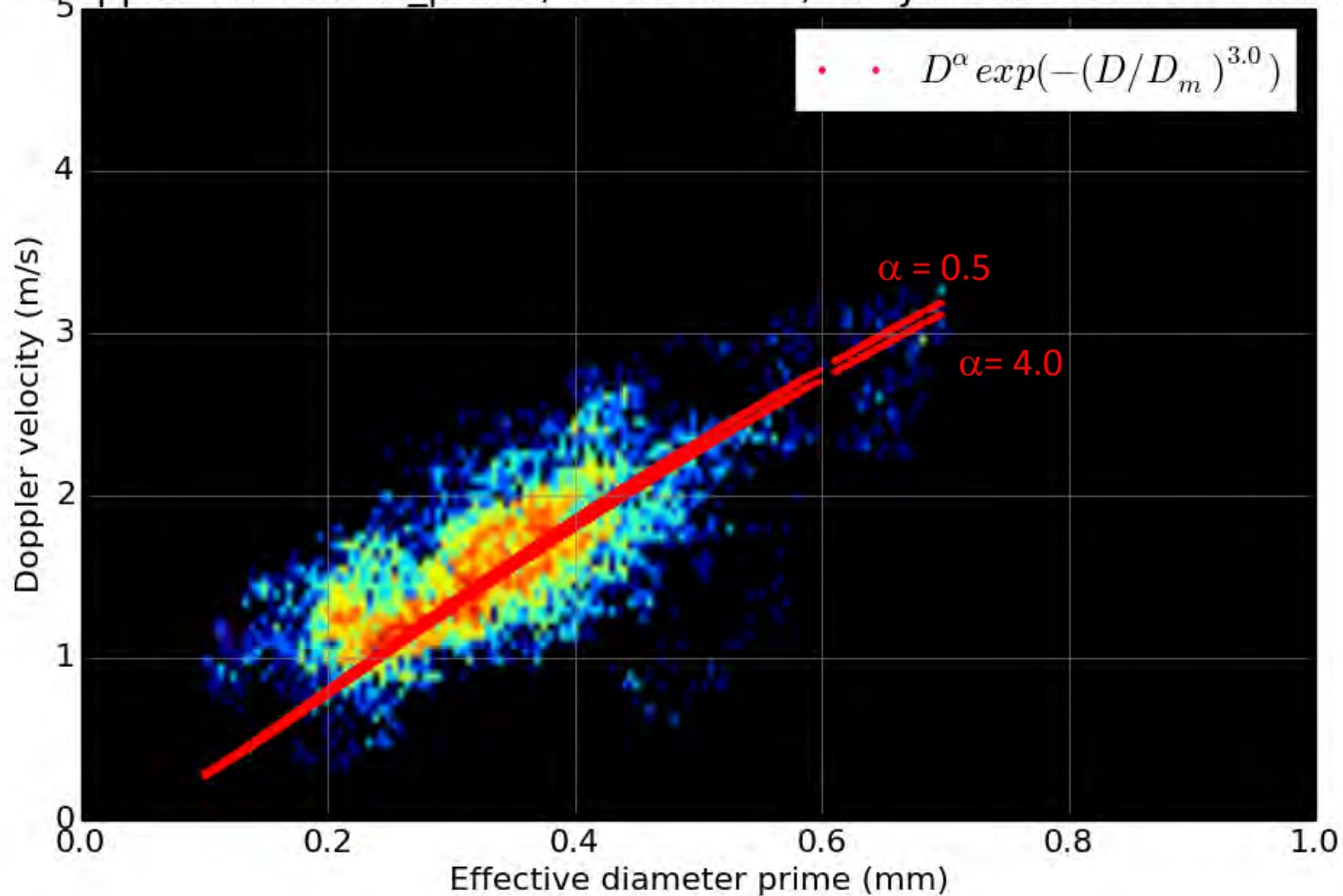
And the Beard Number, X expressed in terms of particle area, volume and density along with the acceleration of gravity, air density, particle diameter and the dynamic viscosity:

$$X = \frac{2 \cdot \text{volume} \cdot \rho_{particle} \cdot \rho_{air} \cdot g \cdot D^2}{\text{area} \cdot \eta^2}$$

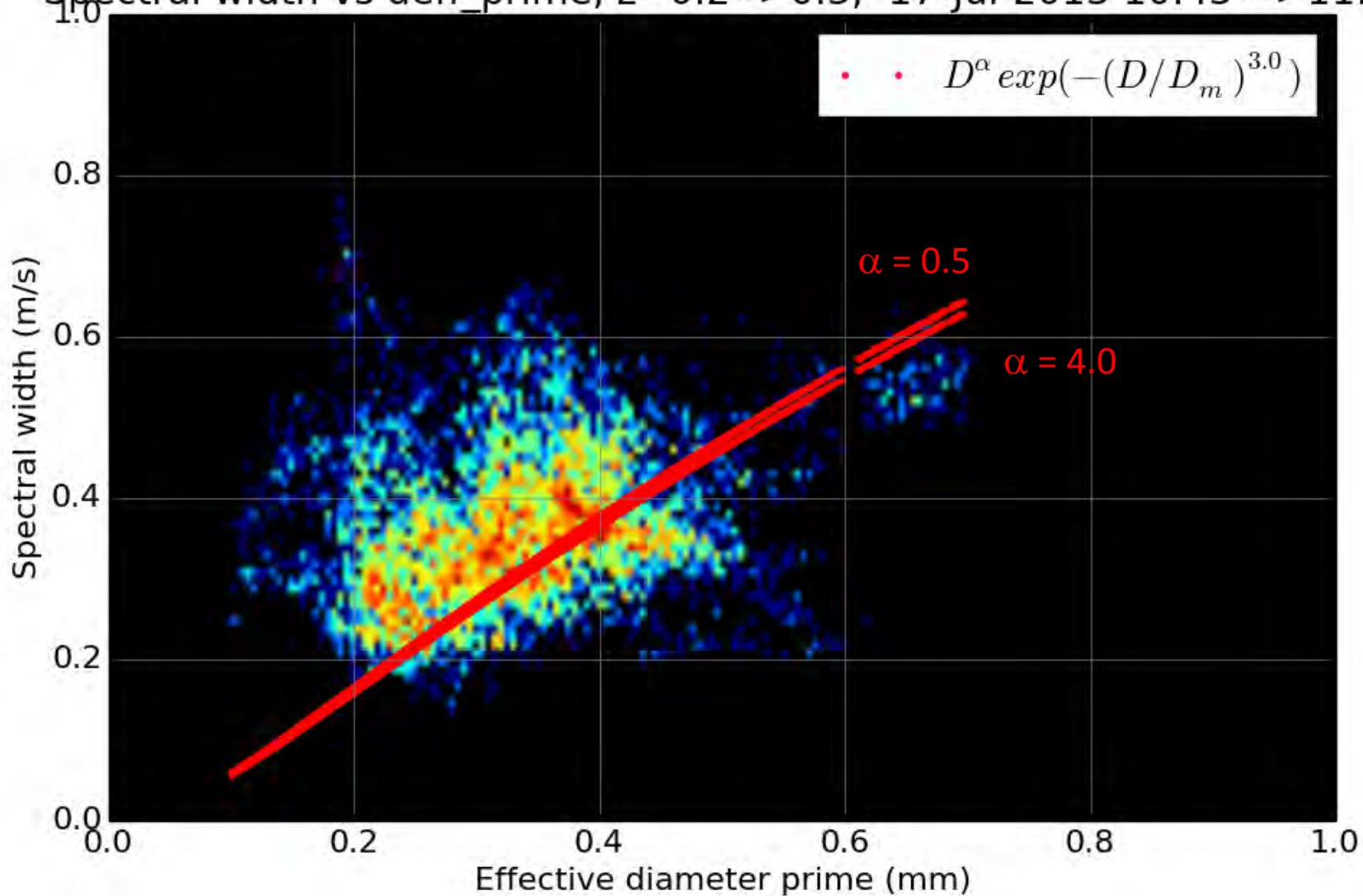
precip rate, dist($\alpha_w = 1.0, \gamma_w = 3.0, \alpha_i = 1.0, \gamma_i = 1.0$) 17-Jul-2013



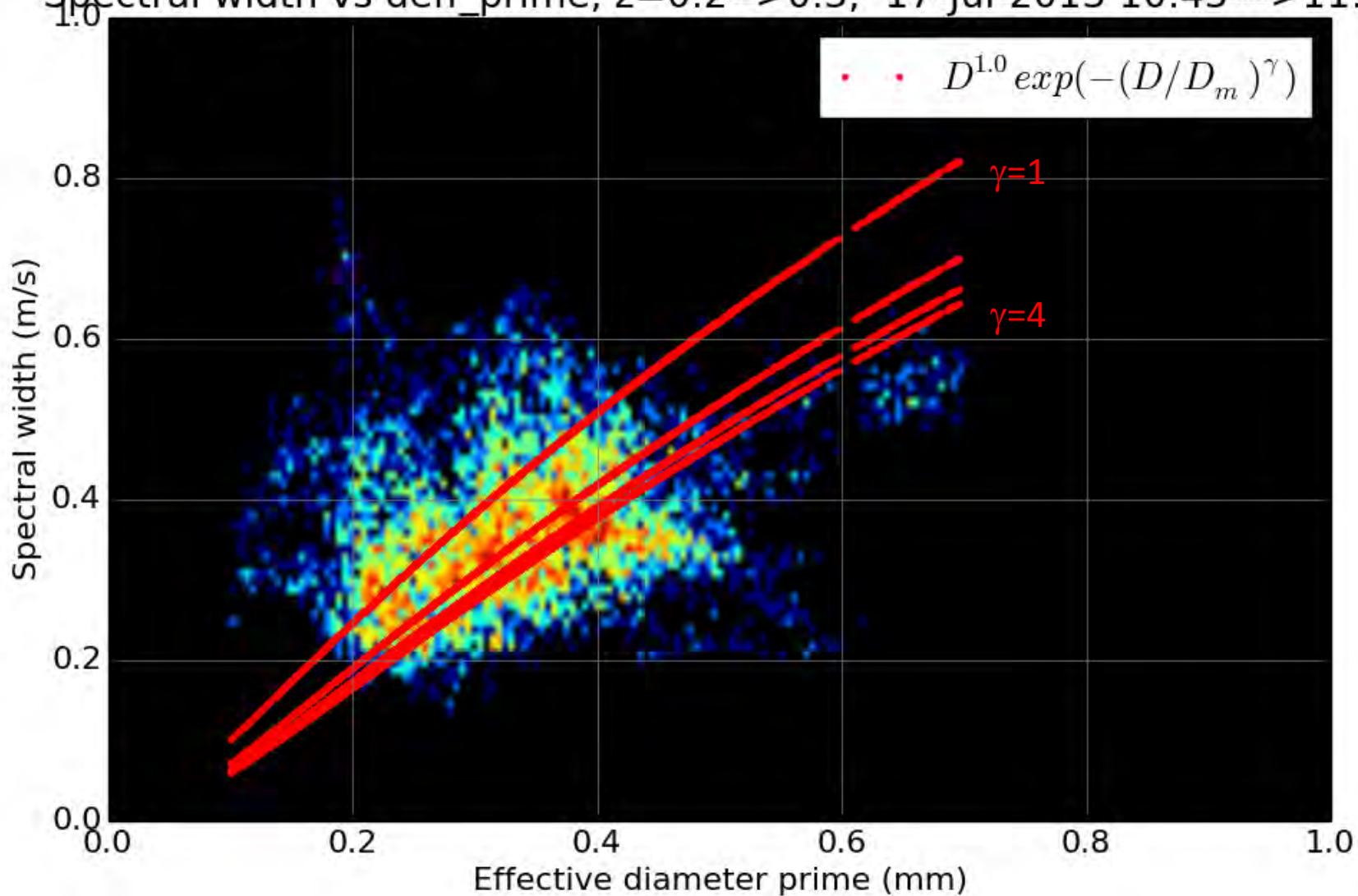
Doppler vel vs deff_prime, z=0.2-->0.3, 17-Jul-2013 10:45--->11:15



Spectral width vs deff_prime, z=0.2-->0.3, 17-Jul-2013 10:45--->11:15

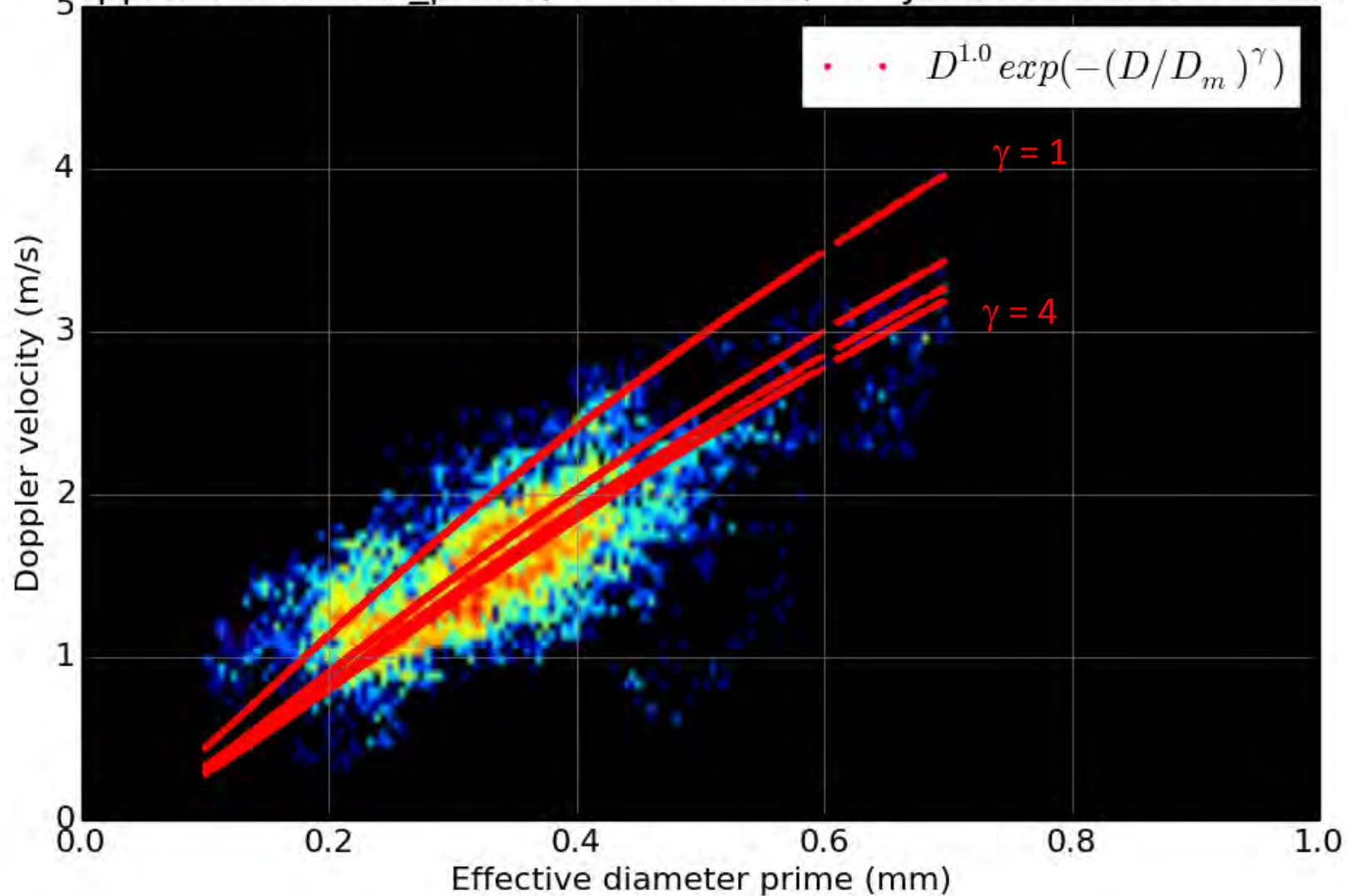


Spectral width vs deff_prime, z=0.2-->0.3, 17-Jul-2013 10:45--->11:15

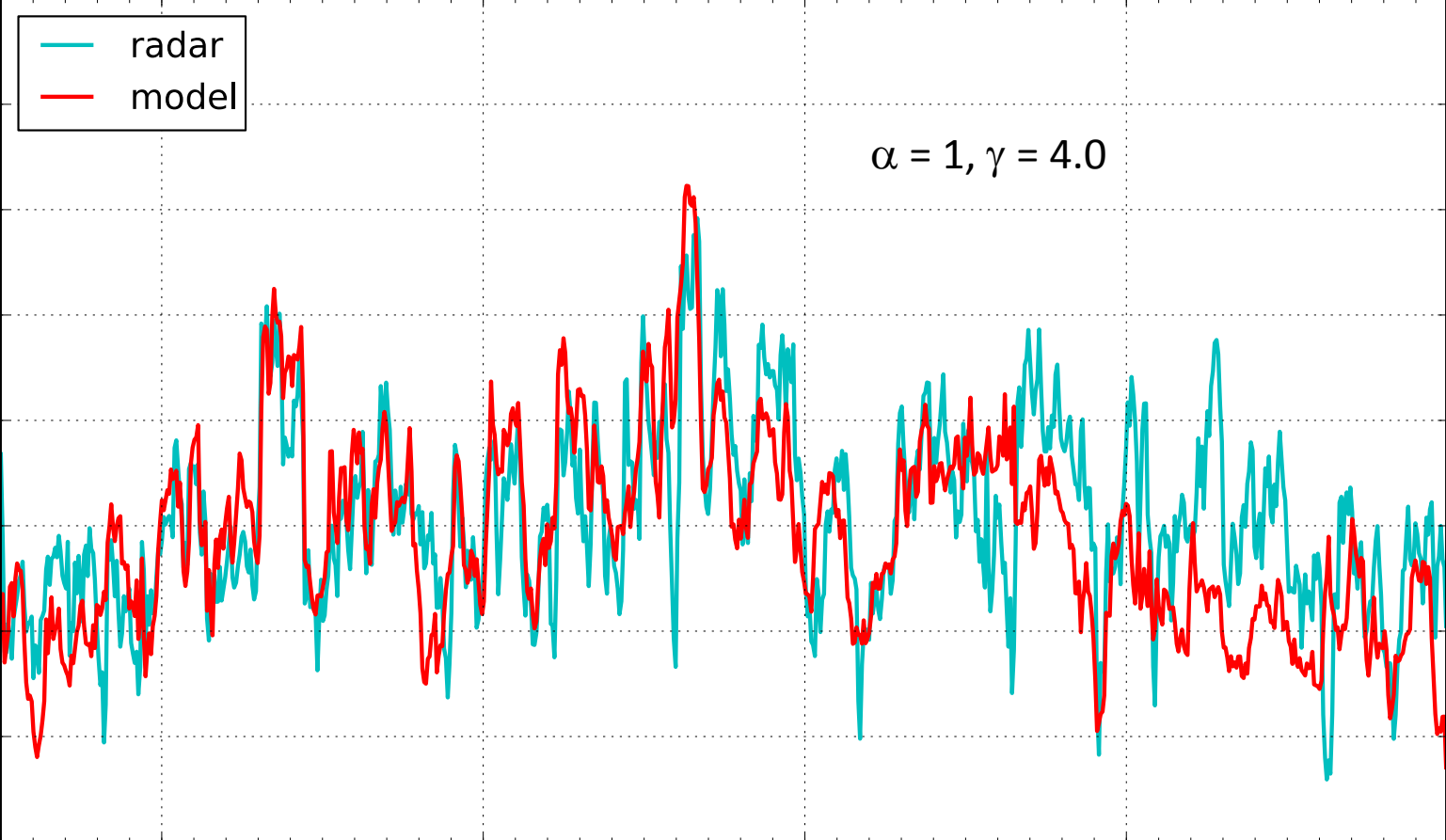


Gamma = 1, 2, 3, 4

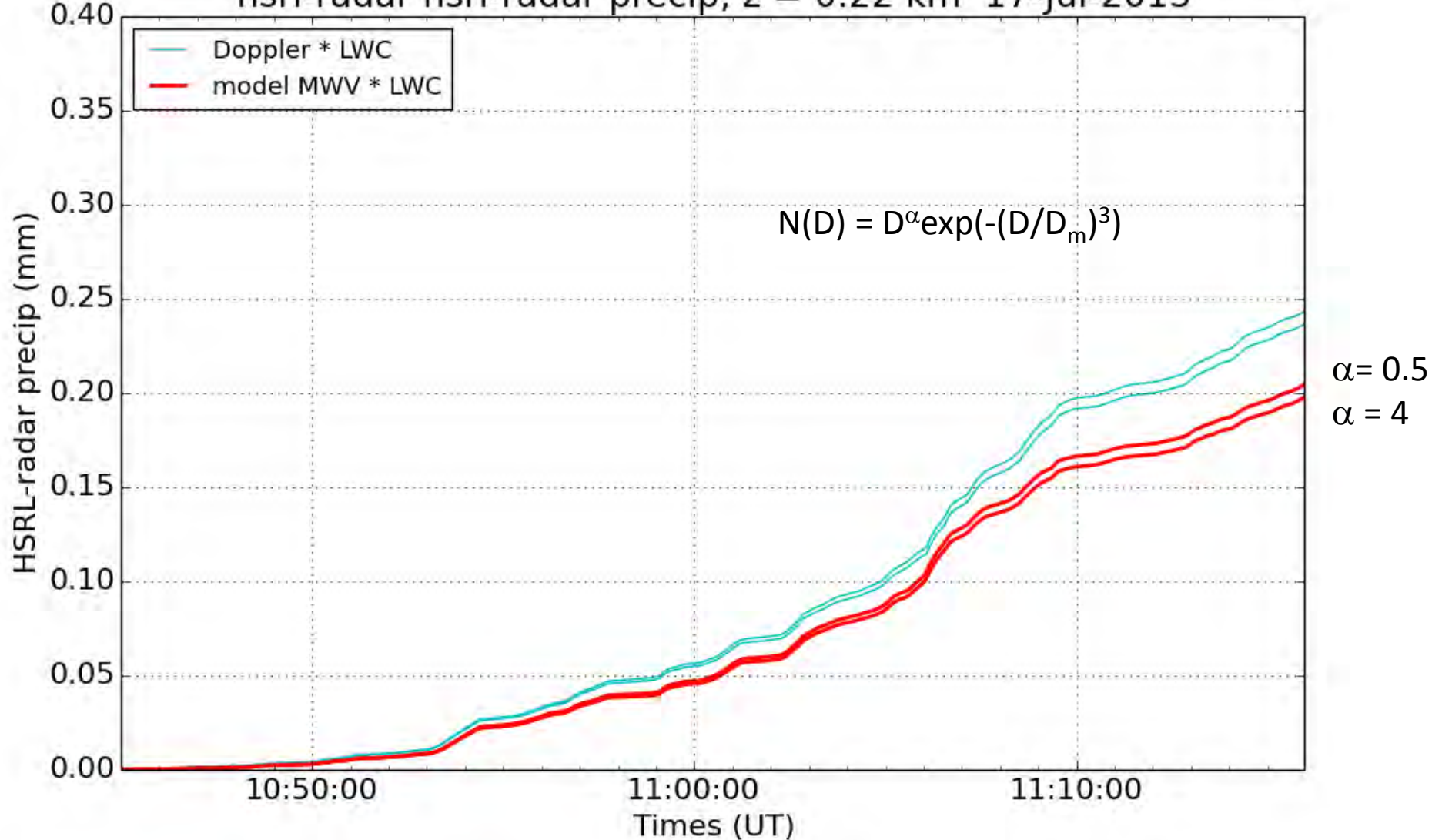
Doppler vel vs deff_prime, z=0.2-->0.3, 17-Jul-2013 10:45--->11:15



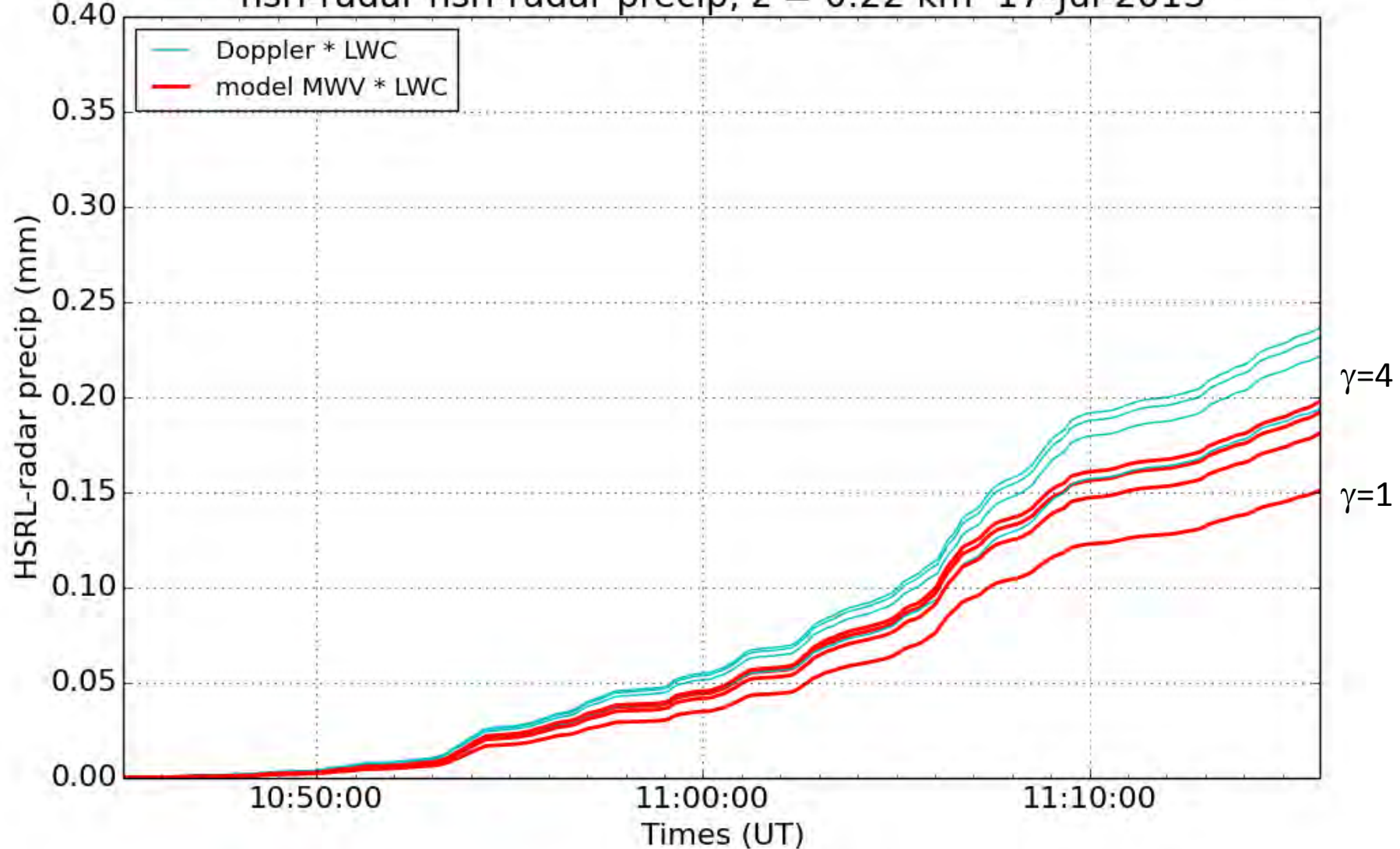
ar-model radar weighted velocity, z= 0.22 km 17-Jul-2013



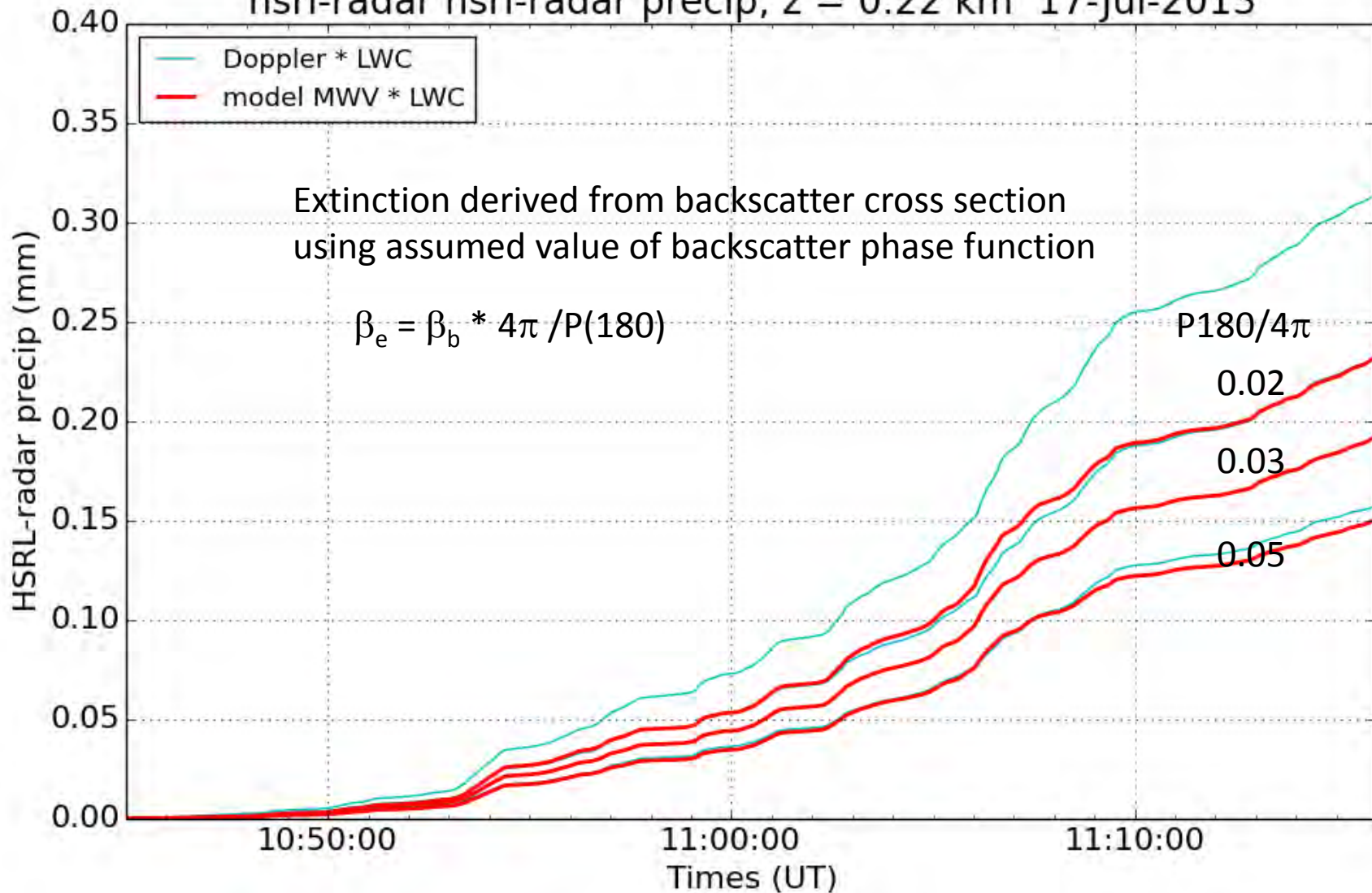
hsrl-radar hsrl-radar precip, z = 0.22 km 17-Jul-2013



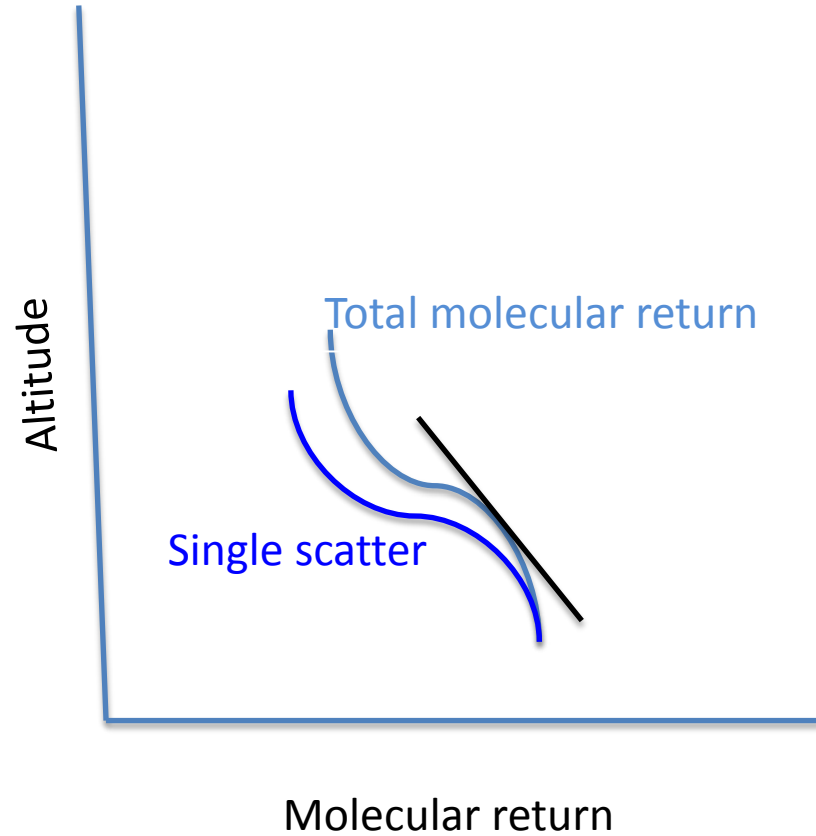
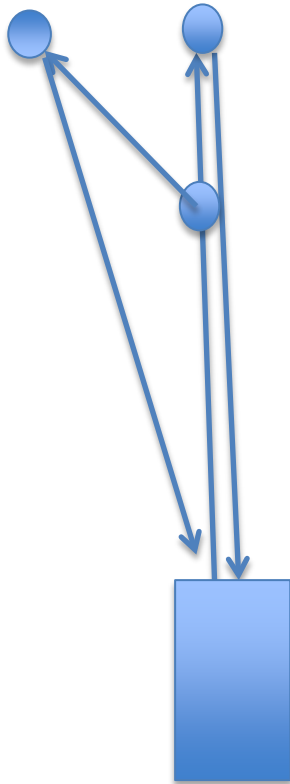
hsrl-radar hsrl-radar precip, z = 0.22 km 17-Jul-2013



hsrl-radar hsrl-radar precip, z = 0.22 km 17-Jul-2013



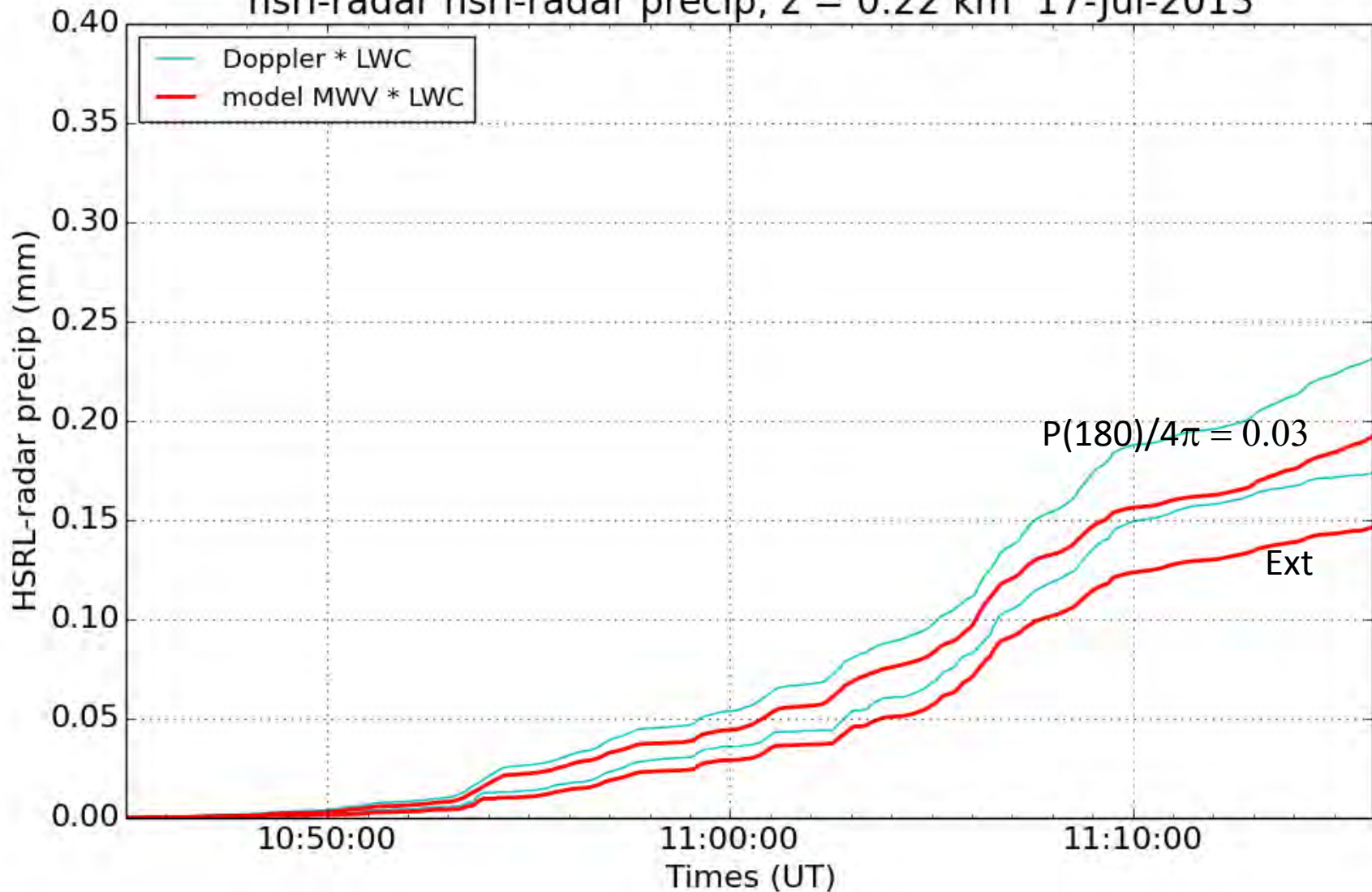
Extinction is derived from the slope of the molecular return
Multiple scattering reduces apparent extinction



Diffraction peak $\theta = \lambda/D \sim 0.5 \text{ micron}/500 \text{ micron} = 1 \text{ mr}$

Extinction derived from $p(180)/4p$ vs direct extinction measurement

hsrl-radar hsrl-radar precip, $z = 0.22$ km 17-Jul-2013



$$\text{Lidar ratio} = 4\pi/P(180)$$

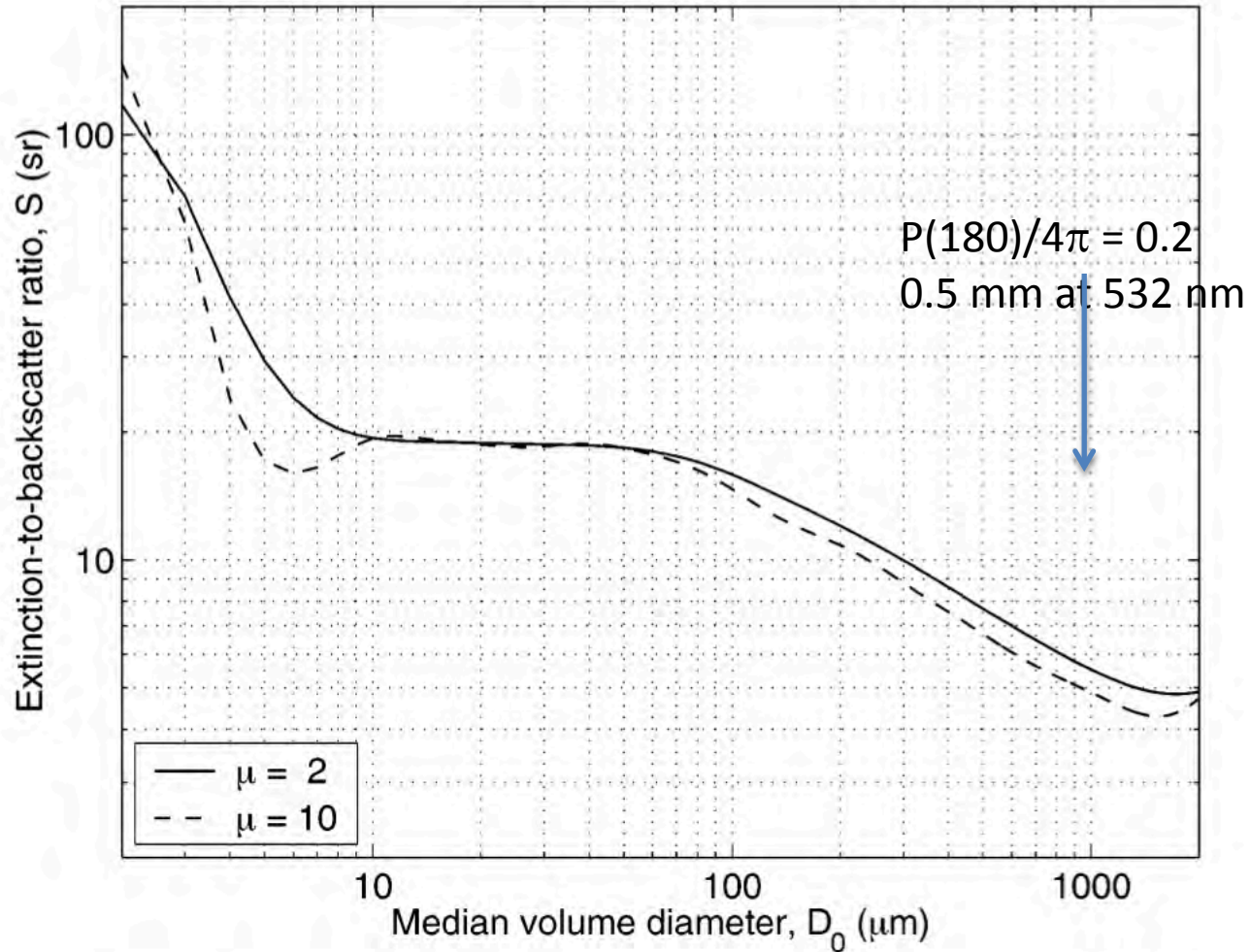
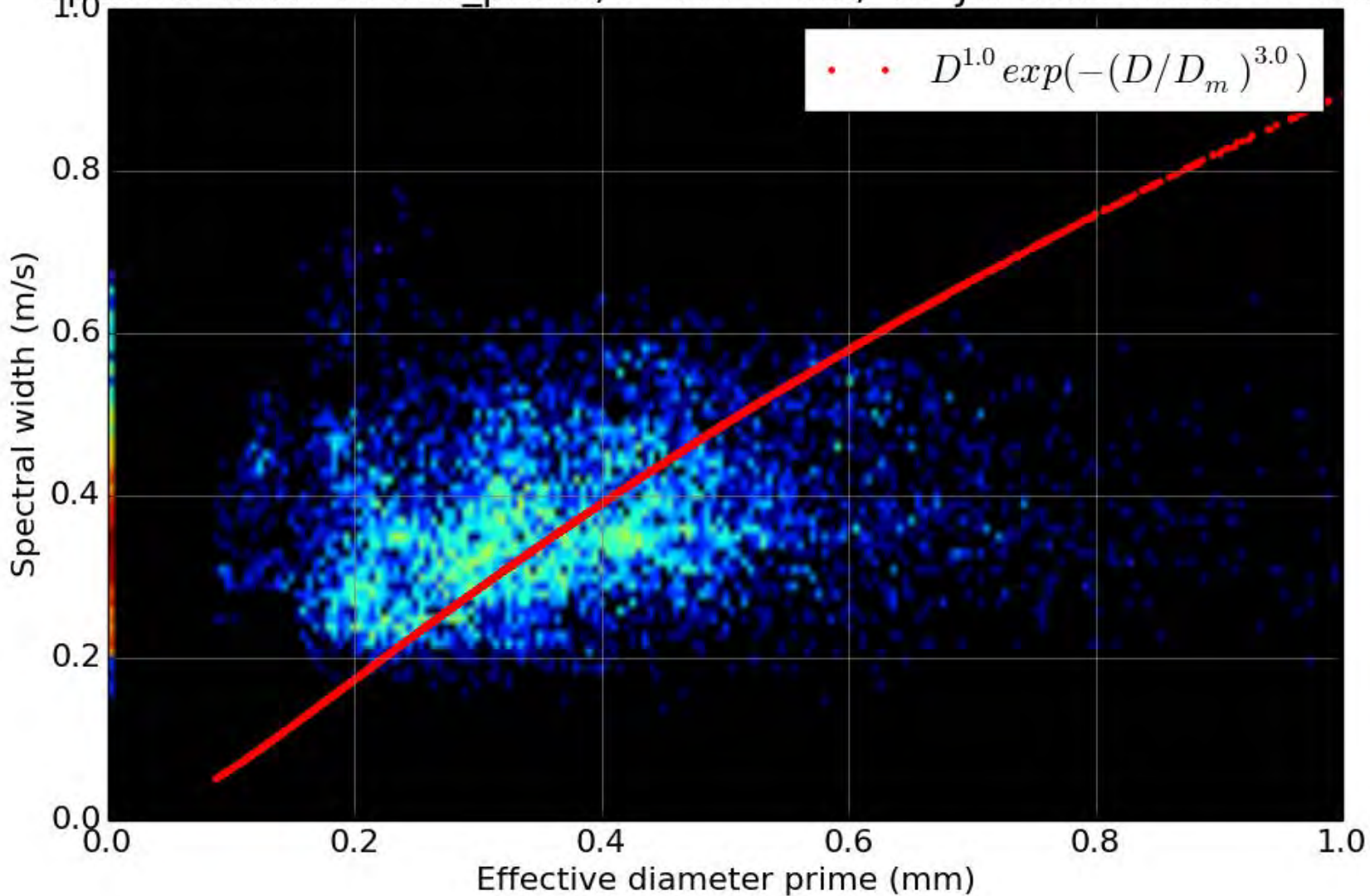
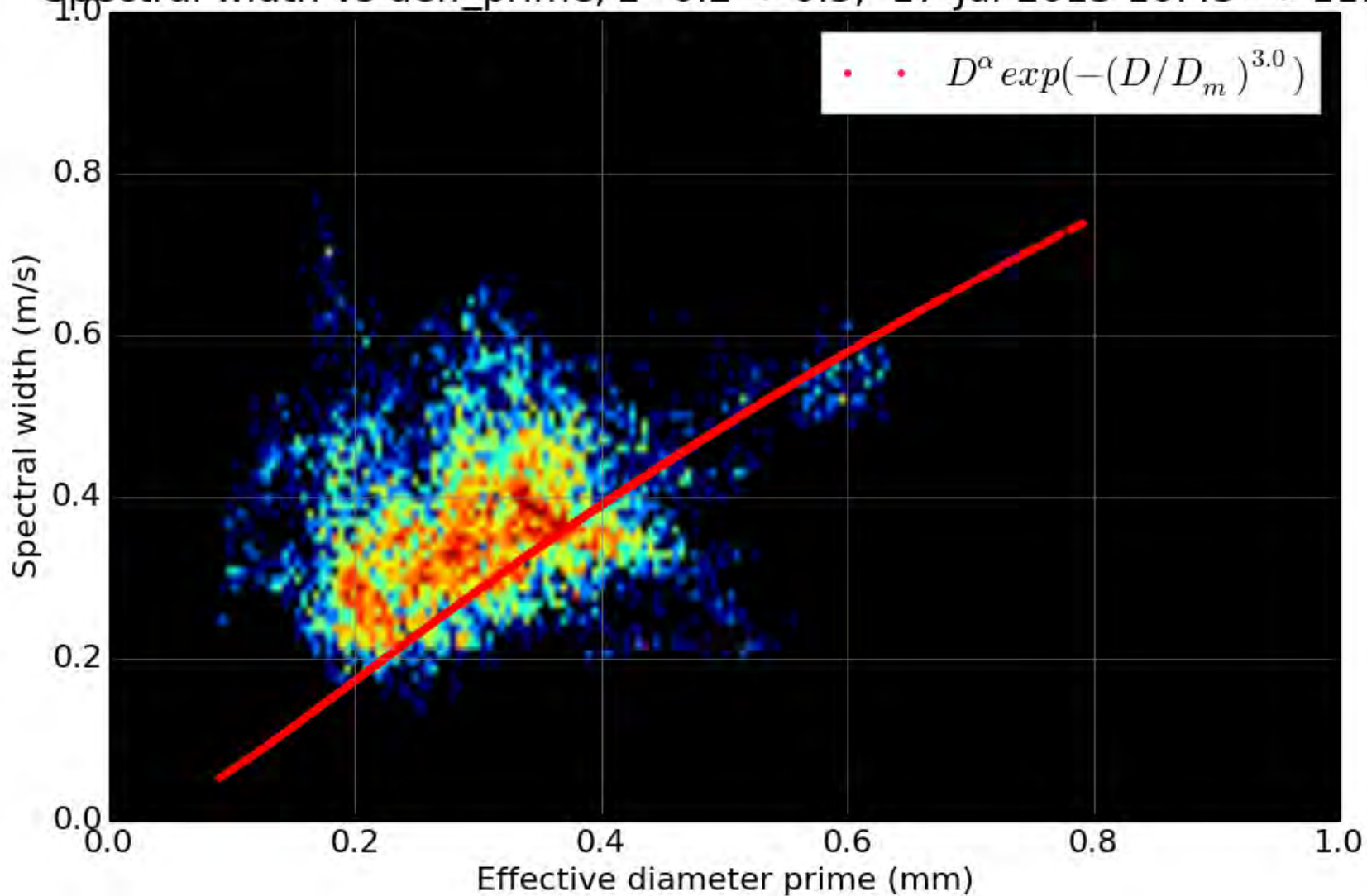


FIG. 3. Theoretical lidar ratio, S , at 905 nm as a function of median volume diameter for gamma distributions of droplet sizes with two different values of μ .

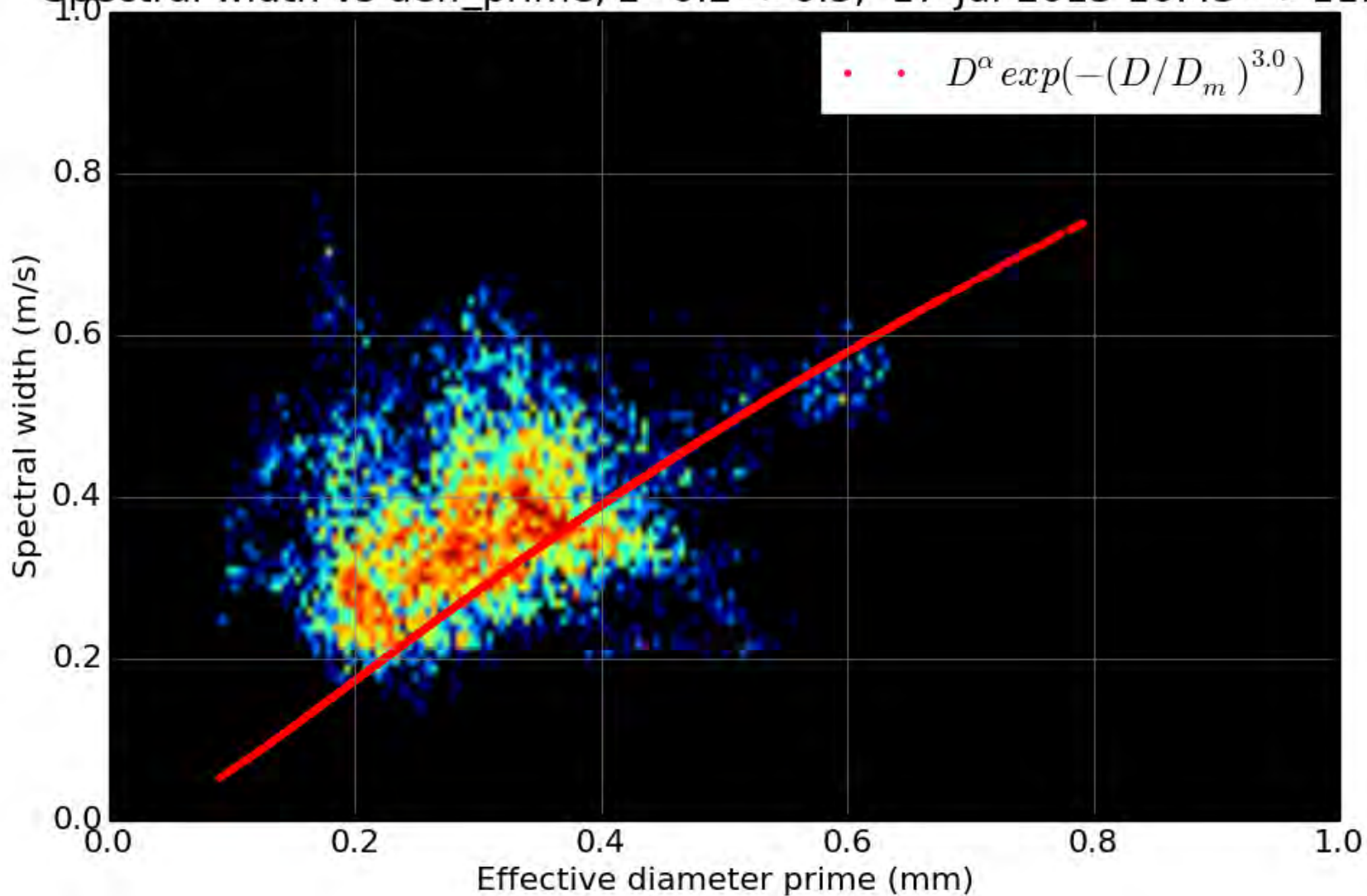
Spectral width vs deff_prime, z=0.2-->0.3, 17-Jul-2013 10:45--->11:15



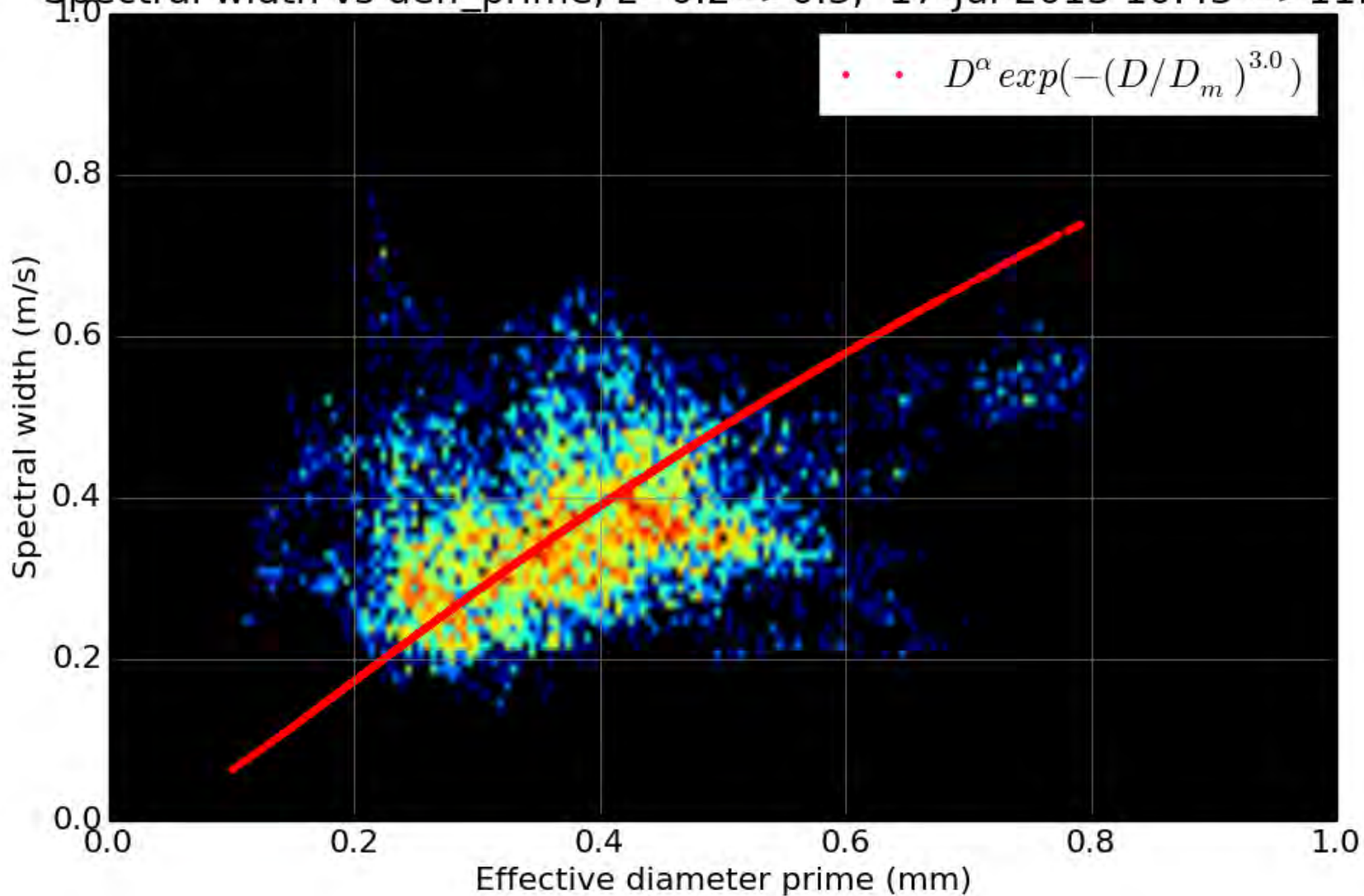
Spectral width vs deff_prime, z=0.2-->0.3, 17-Jul-2013 10:45--->11:15



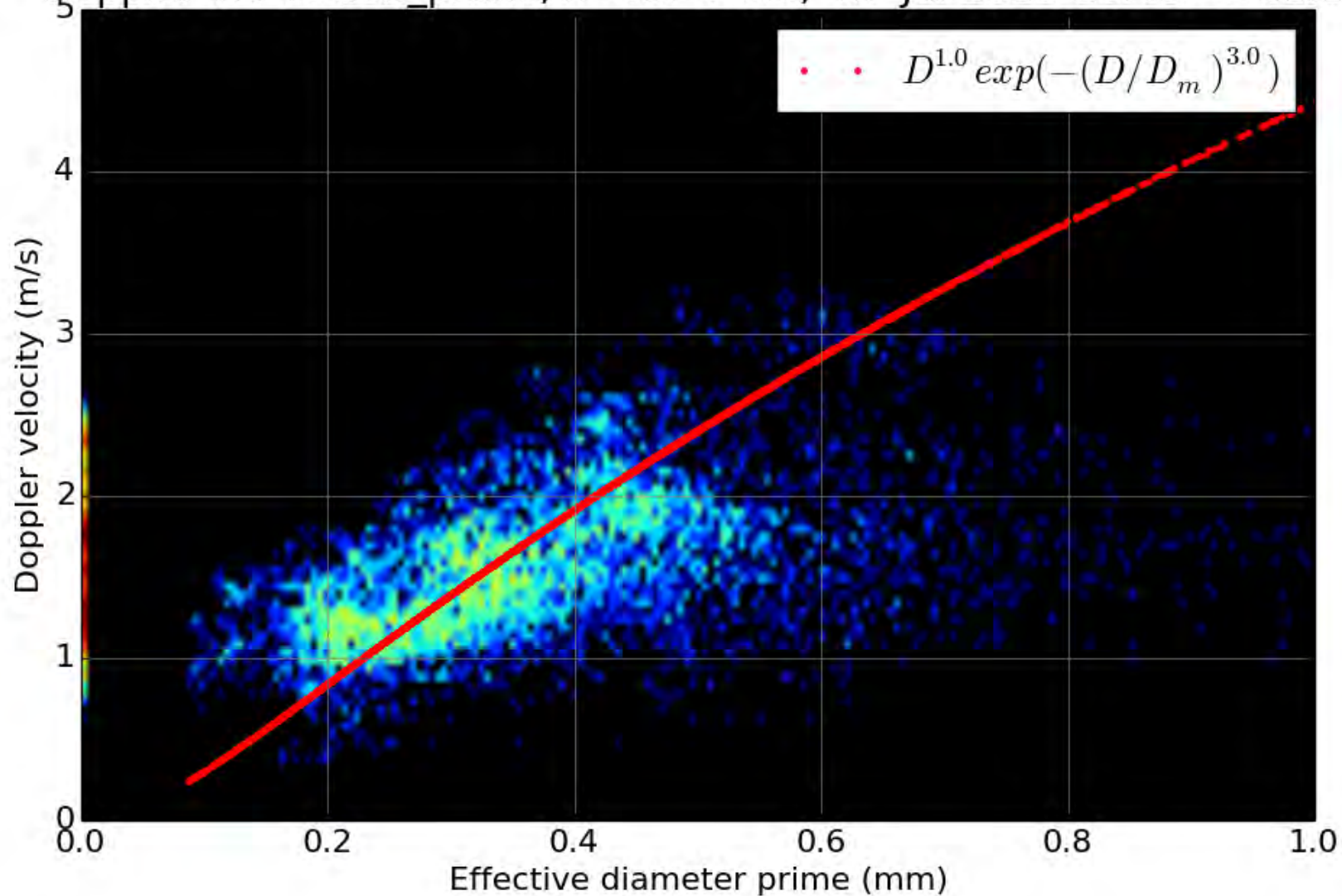
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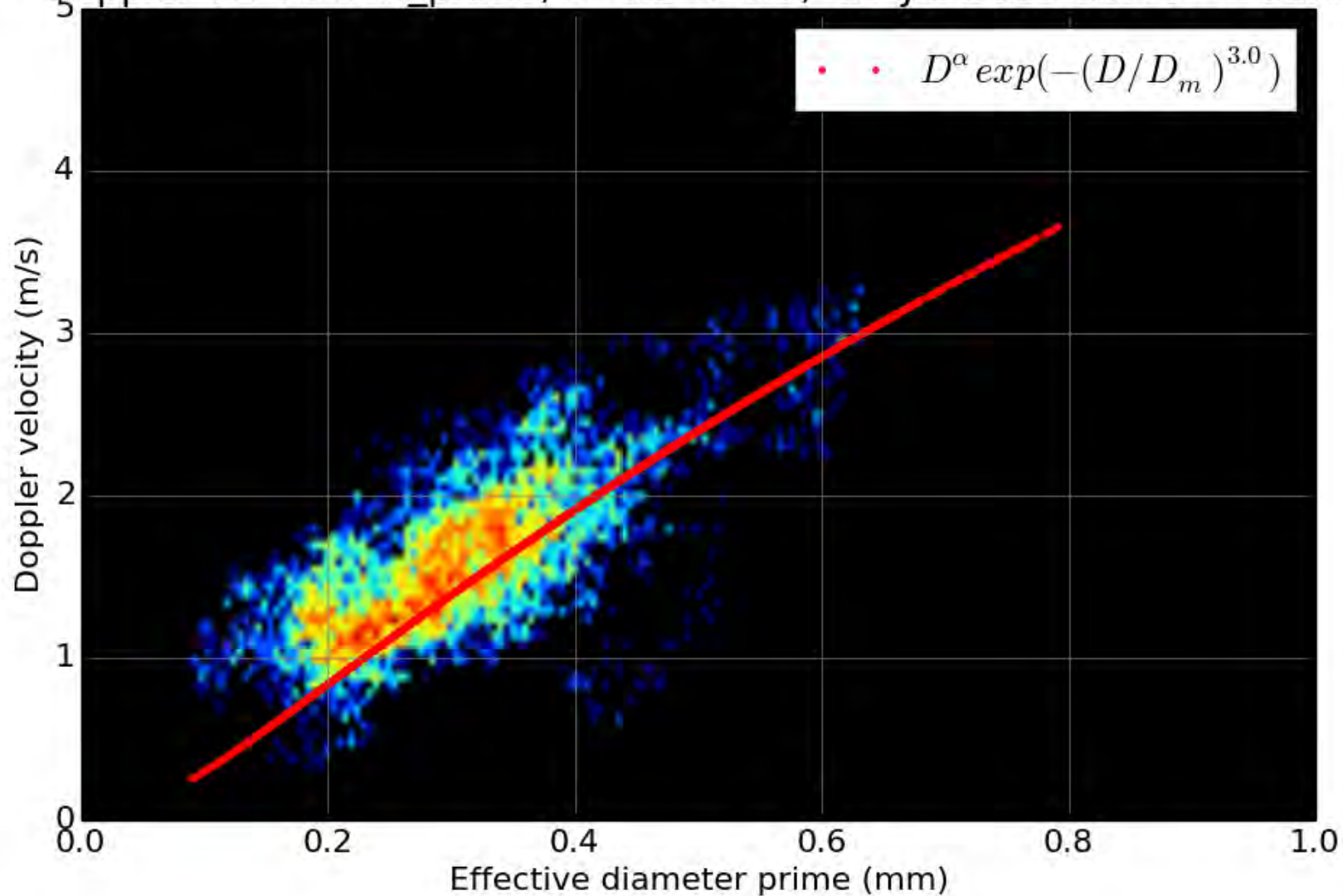
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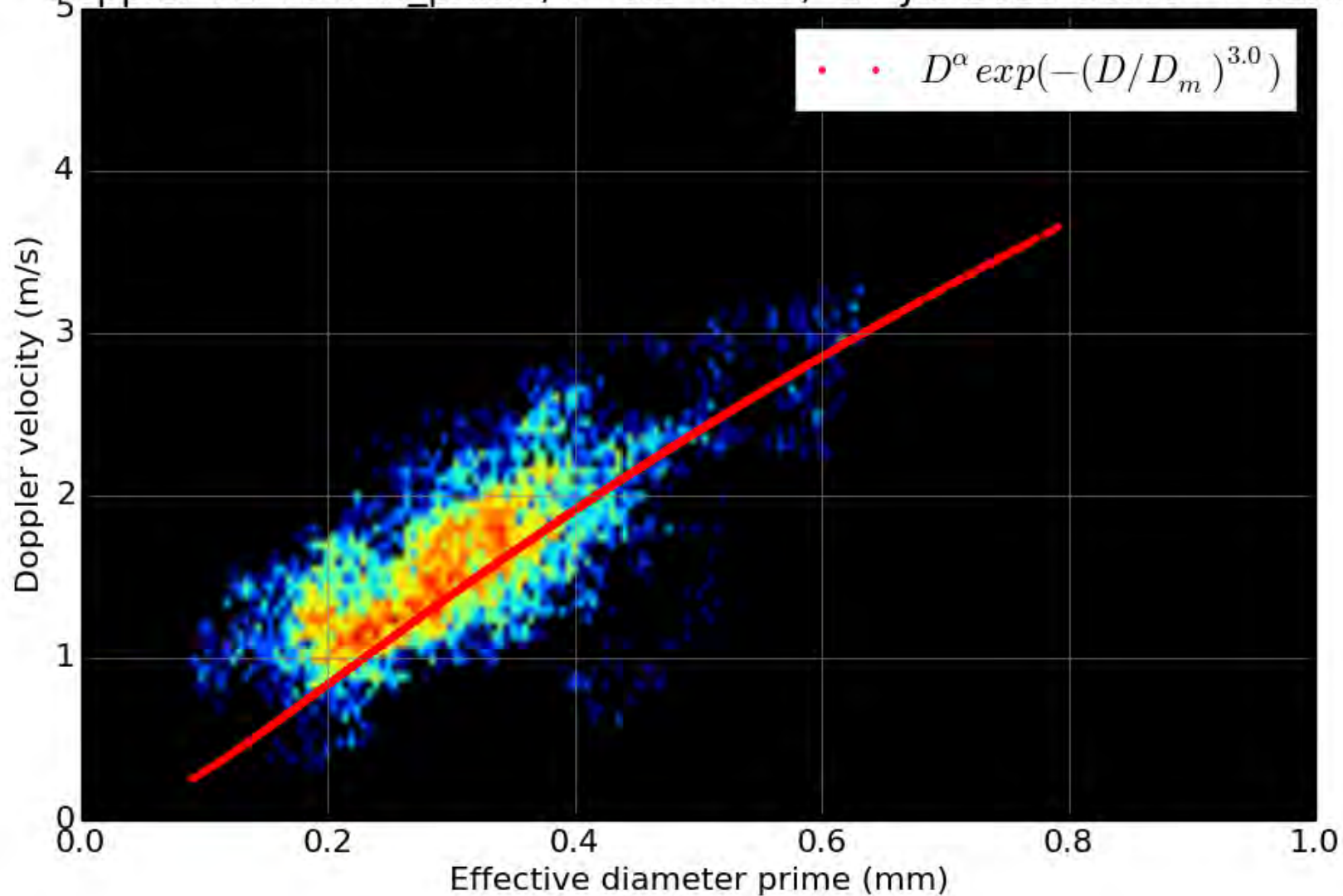
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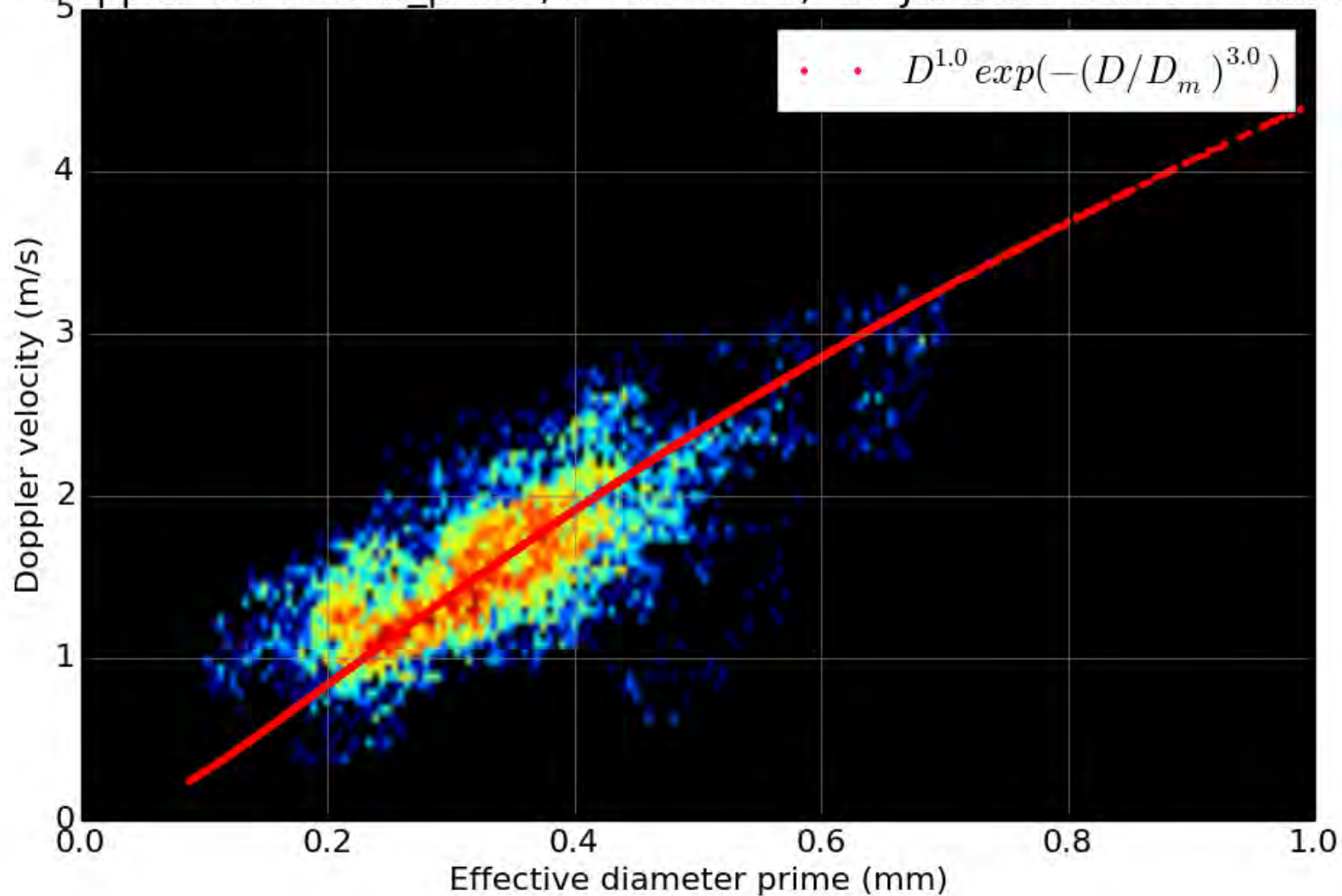
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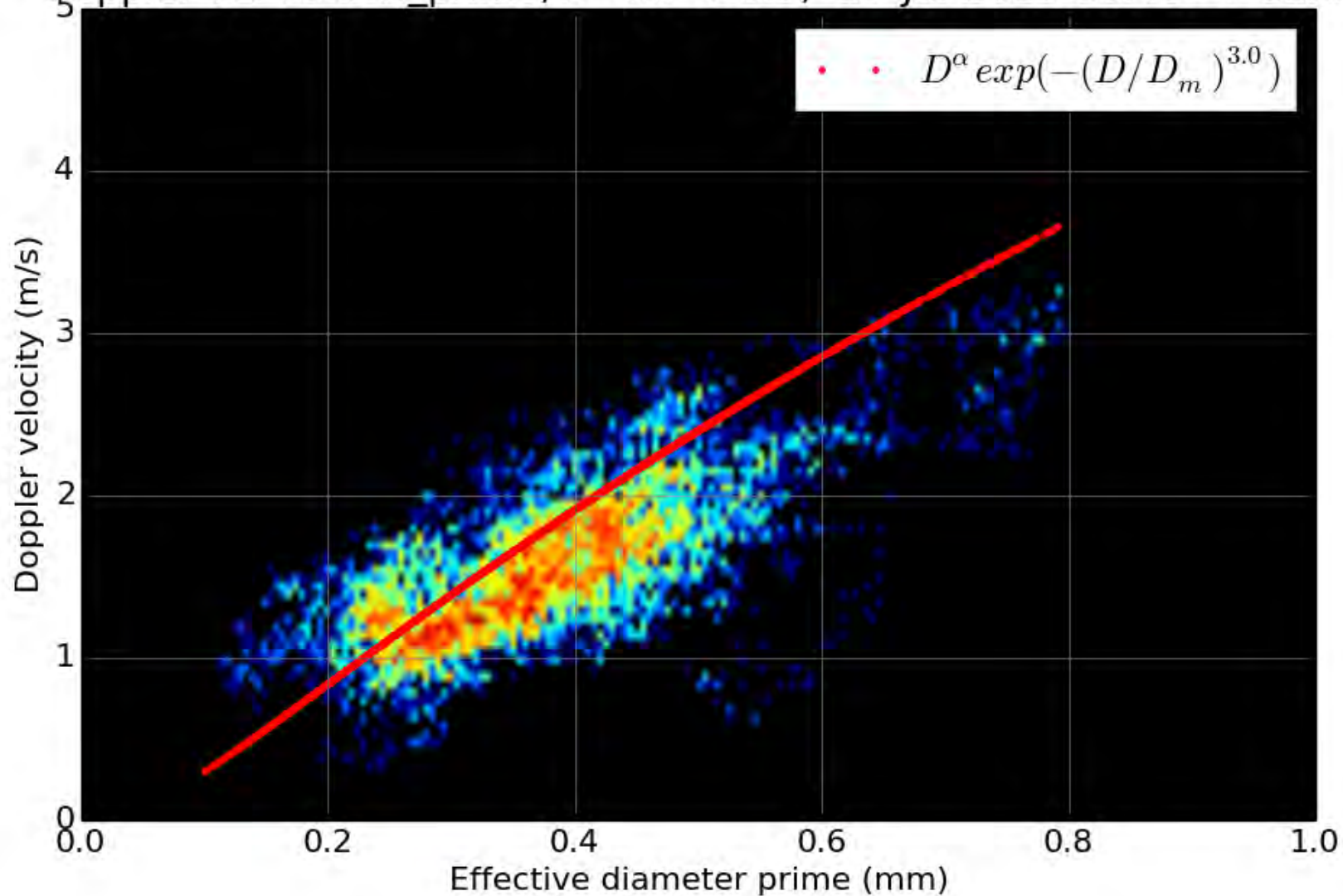
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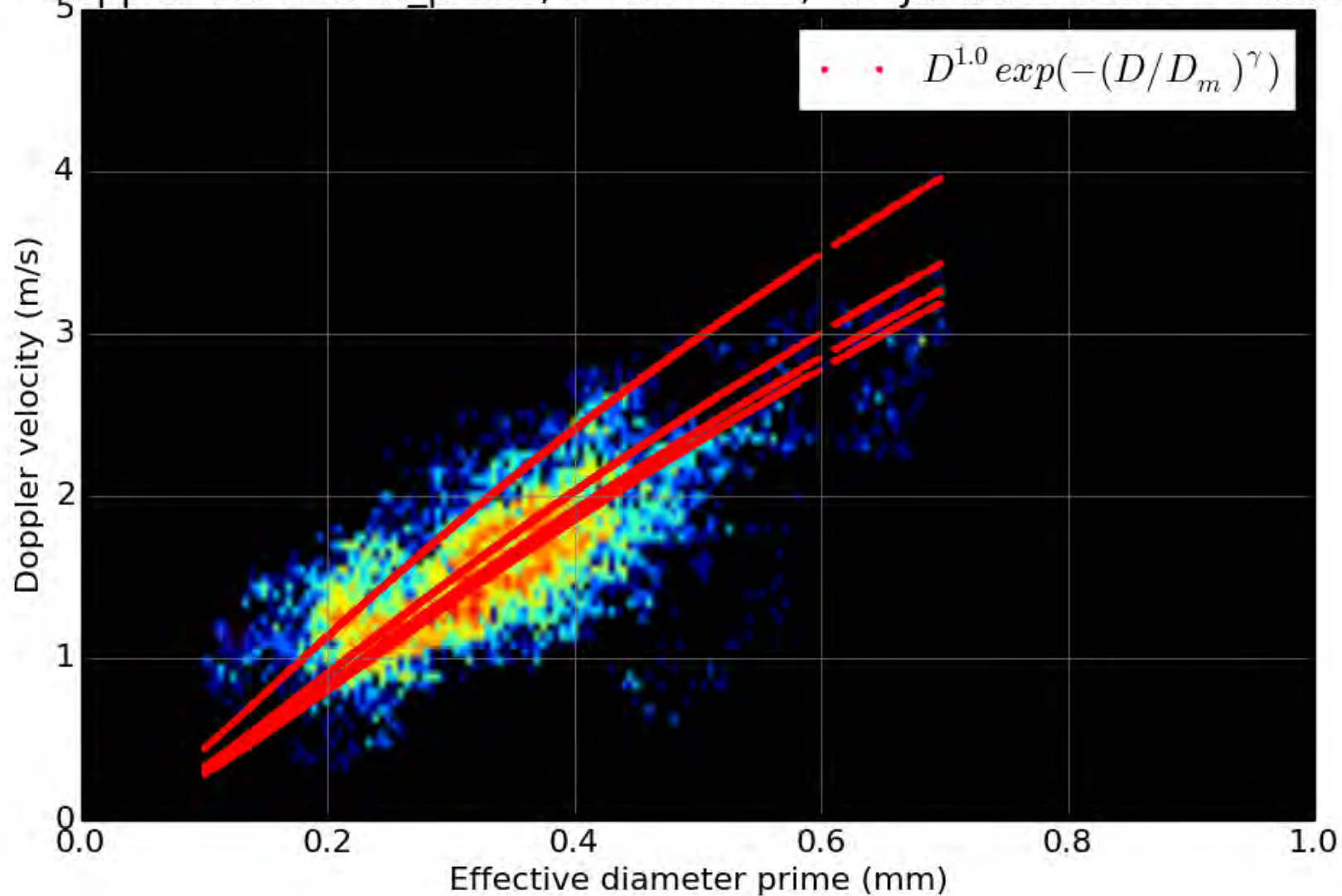
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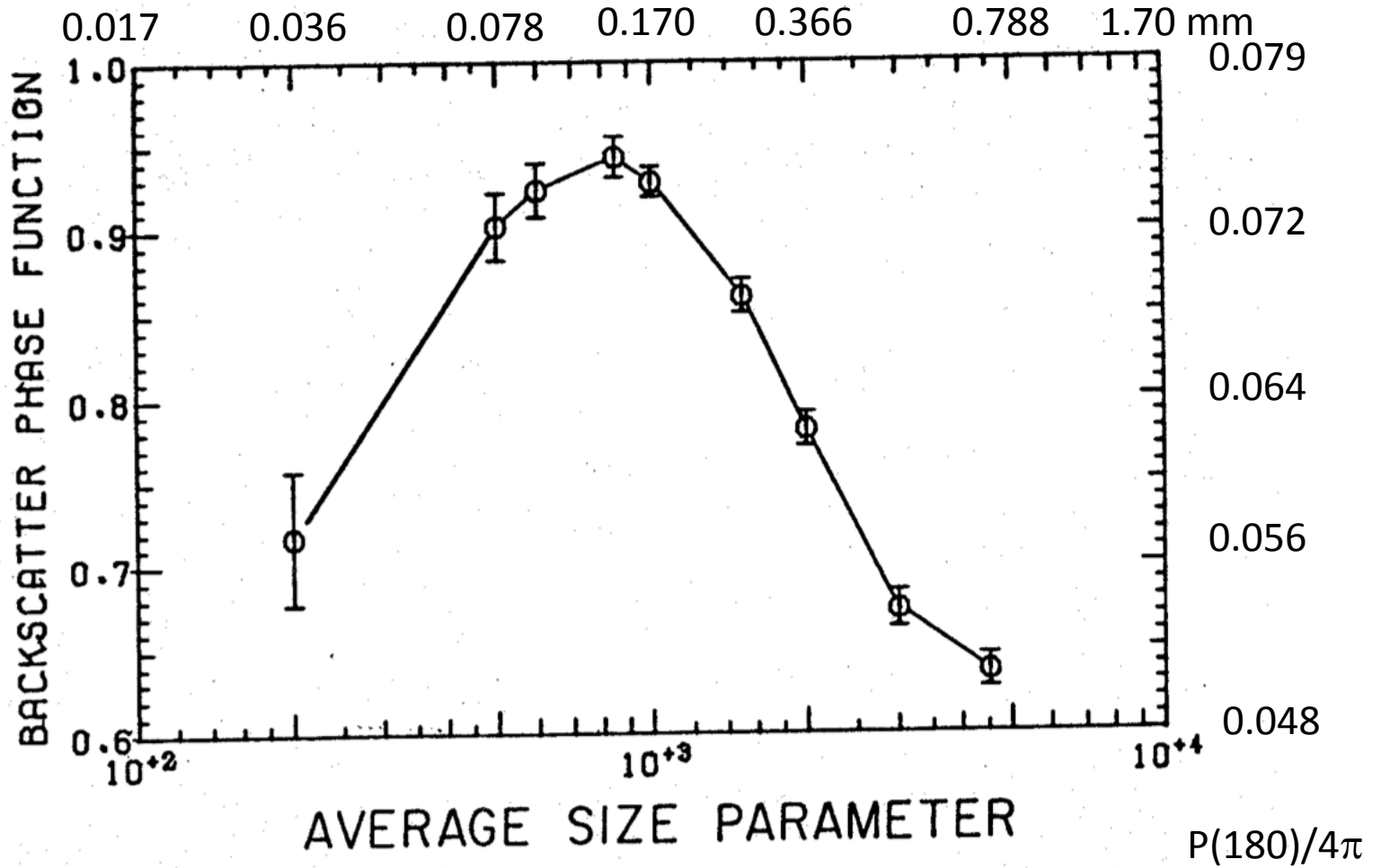
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Doppler vel vs deff_prime, z=0.2-->0.3, 17-Jul-2013 10:45--->11:15



Diameter
 $\lambda = 532\text{nm}$



Backscatter phase function for water drops averaged over size parameter intervals of 0.14. (Shiple 1978)